

A NEW TECHNIQUE TO FIND LAPLACE TRANSFORM

AMIT GAUR and GARIMA AGARWAL

Department of Mathematics and Statistics Manipal University Jaipur, Jaipur Rajasthan, India-303007 E-mail: gaur84amit@gmail.com

Abstract

Recently a new integral transform named β -Laplace integral transform has been introduced which is a new form of generalization of classical Laplace transform. This new generalized transform has many characteristic properties over Laplace transform. In this work, we purpose a new technique to find Laplace transform by applying the β -Laplace integral transform associated with Leibnitz rule which is simple and efficient than the direct Laplace transform technique.

1. Introduction

Integral transforms always been a very useful mathematical tool to solve many kind of differential equation, partial differential equation, integral equations. We have a wide range of integral transform such as Laplace, Fourier, Hankel, Mellin, Radon, Gabor, Hilbert, Weiestrauss, Abel, Sumudu, etc. They all have a wide range of applications in the fields of Physical science, Mathematics, Statistics, and engineering [2, 9, 11].

Laplace integral transform is one of the oldest and famous integral transforms purposed by P. S. Laplace in its most celebrated work 'Théorie analytique des probabilités' [8]. Laplace transform has a wide range of applications in many fields of science and engineering [1, 3, 10, 12].

Recently a new form of generalization of Laplace Transform has been introduced and named β -Laplace integral transform which has many characteristic properties over Laplace transform [4, 5, 6].

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One more beautiful aspect of this new β -Laplace integral transform is that it is generalization of most of the recently introduced transforms, for instance, Sumudu, Kamal, Natural, Polynomial, Tarig, Elzaki, Aboodh, Laplace-Carson (Mahgoub), Mohand, Sawi, Sadik integral transform [6].

In this paper, we establish a connection between β -Laplace integral transform and Laplace integral transform via Leibnitz rule to find Laplace integral transform of a function. Here we purpose a new efficient and simple technique to find Laplace transform by using this new generalized β -Laplace transform via Leibnitz rule. We illustrate this technique by some examples.

2. Definitions and Prepositions

(i) Laplace Integral Transform

Laplace integral transform of a suitable function $\psi(t)$, $t \ge 0$ is defined by

$$L\{\psi\}_{(p)} = \int_0^\infty e^{-pu} \psi(u) du, \ p \in \mathbb{C}, \ \operatorname{Re}(p) > 0.$$
(1)

Where suitable function means a function for which improper integral of right side converges.

(ii) β-Laplace Integral Transform [4]

β-Laplace integral transform of a suitable function $\psi(t), t \ge 0$ is defined by

$$L_{\beta}\{\psi\}_{(p)} = \int_0^\infty \beta^{-pu} \psi(u) du, \ \beta > 1, \ p \in \mathbb{C}, \ \operatorname{Re}(p) > 0.$$
⁽²⁾

(iii) Leibnitz Integral Rule [7]

Leibnitz rule (Differentiation under integral sign) for an integral in the form $\int_{\phi_1(x)}^{\phi_2(x)} \psi(x, y) dy$, states that:

Let $\psi(x, y)$ be a continuous function with continuous partial derivative $\frac{\partial \psi(x, y)}{\partial x}$ in some region of the (x, y)-plane which include region $\phi_1(x) \le y \le \phi_2(x), x_0 \le x \le x_1$, where both the functions $\phi_1(x)$, and $\phi_2(x)$

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are continuous and have continuous derivatives for $x_0 \le x \le x_1$. Then, for $x_0 \le x \le x_1$,

$$\frac{d}{dx} \left\{ \int_{\phi_1(x)}^{\phi_2(x)} \psi(x, y) dy \right\} = \int_{\phi_1(x)}^{\phi_2(x)} \frac{\partial}{\partial x} \{ \psi(x, y) \} dy + \psi(x, \phi_1(x)) \cdot \frac{d}{dx} \{ \phi_1(x) \} - \psi(x, \phi_2(x)) \cdot \frac{d}{dx} \{ \phi_2(x) \}.$$
(3)

Preposition 1 ([4]). Let $\psi(t)$, $t \ge 0$ be a suitable function then

$$\lim_{\beta \to e} L_{\beta} \{\psi\}_{(p)} = L\{\psi\}_{(p)}.$$
(4)

Preposition 2 ([4]). Let $\psi(t), t \ge 0$ be a suitable function then

$$\lim_{\beta \to e} L_{\beta} \{\psi\}_{(p)} = 0.$$
(5)

3. Main Results

In this section we introduce a new way to find Laplace transform with the help of applying Leibnitz rule on β -Laplace integral transform.

Let $\psi(t)$ be the function under consideration of finding Laplace integral transform, and we start from the definition of β -Laplace integral transform

$$L_{\beta}\{\psi\}_{(p)} = \int_0^\infty \beta^{-pu} \psi(u) du.$$

Taking partial derivative both sides

$$\frac{\partial}{\partial\beta} \left\{ L_{\beta} \{ \psi \}_{(p)} \right\} = \frac{\partial}{\partial\beta} \int_{0}^{\infty} \beta^{-pu} \psi(u) du.$$

After applying Leibnitz rule, we get

$$\frac{\partial}{\partial\beta} \left\{ L_{\beta} \{\psi\}_{(p)} \right\} = \int_{0}^{\infty} \left(-\frac{pu}{\beta} \right) \beta^{-pu} \psi(u) du.$$

Let find the β -Laplace transform of the right side of the equation, plug into equation

$$\frac{\partial}{\partial \beta} \{ L_{\beta} \{ \psi \}_{(p)} \} = -\frac{p}{\beta} G(\beta, p).$$

Where $\int_0^\infty u\beta^{-pt}\psi(u)du = G(\beta, p).$

Taking integration both sides with respect to $\boldsymbol{\beta}$

$$L_{\beta}\{\psi\}_{(p)} = -\int \frac{p}{\beta} G(\beta, p) d\beta.$$

 $L_{\beta}\{\psi\}_{(p)} = \overline{G}(\beta, p) + c, c$ is an integral constant

Where
$$-\int \frac{p}{\beta} G(\beta, p) d\beta = \overline{G}(\beta, p).$$

By preposition (2)

$$\lim_{\beta\to\infty}L_{\beta}\{\psi\}_{(p)}=0.$$

After taking Limit $\beta \to \infty$

$$\lim_{\beta \to \infty} L_{\beta} \{\psi\}_{(p)} = \lim_{\beta \to \infty} (\overline{G}(\beta, p) + c)$$
$$\Rightarrow c = -\lim_{\beta \to \infty} \overline{G}(\beta, p).$$

After plugging the value of c in equation

$$L_{\beta}\{\psi\}_{(s)} = \overline{G}(\beta, p) - \lim_{\beta \to \infty} \overline{G}(\beta, p).$$

Finally, after taking limit $\beta \rightarrow e$, we obtain our desired Laplace transform

$$\begin{split} \lim_{\beta \to e} L_{\beta} \{\psi\}_{(p)} &= \lim_{\beta \to e} \left(\overline{G}(\beta, p) - \lim_{\beta \to e} \overline{G}(\beta, p)\right) \\ L\{\psi\}_{(p)} &= \lim_{\beta \to e} \overline{G}(\beta, p) - \lim_{\beta \to e} \left(\lim_{\beta \to \infty} \overline{G}(\beta, p)\right) \\ L\{\psi\}_{(p)} &= \lim_{\beta \to e} \int \left(-\frac{p}{\beta}\right) G(\beta, p) d\beta - \lim_{\beta \to e} \left\{\lim_{\beta \to \infty} \int \left(-\frac{p}{\beta}\right) G(\beta, p) d\beta \right\} \\ L\{\psi\}_{(p)} &= \lim_{\beta \to e} \int \left(\int_{0}^{\infty} \left(-\frac{p}{\beta}\right) \beta^{-pu} \psi(u) du\right) d\beta \\ &- \lim_{\beta \to e} \left\{\lim_{\beta \to \infty} \int \left(\int_{0}^{\infty} \left(-\frac{pu}{\beta}\right) \beta^{-pu} \psi(u) du\right) d\beta \right\}. \end{split}$$

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Now We have obtained our desired result.

4. Illustrative Examples

Example 1. The Laplace transform of $\frac{\sin(t)}{t}$, t > 0. By the definition of

 $\beta\text{-Laplace transform}$

$$L_{\beta}\{\psi\}_{(p)} = \int_{0}^{\infty} \beta^{-pu} \psi(u) du$$

Taking $\psi(t) = \frac{\sin t}{t}$, t > 0, we have

$$L_{\beta}\{\psi\}_{(p)} = \int_0^\infty \frac{\beta^{-su} \sin(u)}{u} du.$$

Taking partial derivative with respect to β (Partial derivative under integral sign) We obtain

$$\begin{split} \frac{\partial}{\partial\beta} (L_{\beta}\{\psi\}_{(p)}) &= \frac{\partial}{\partial\beta} \int_{0}^{\infty} \frac{\beta^{-pu} \sin(u)}{u} du = \int_{0}^{\infty} \left(-\frac{pu}{\beta}\right) \frac{\beta^{-pu} \sin(u)}{u} du \\ &= -\frac{p}{\beta} \int_{0}^{\infty} \beta^{-pu} \sin(u) du \\ &= -\frac{p}{\beta} \left(\frac{1}{(p \ln \beta)^{2} + 1}\right). \end{split}$$

Integrating with respect to β

$$L_{\beta}\{\psi\}_{(p)} = \int \left(-\frac{p}{\beta}\left(\frac{1}{\left(p\ln\beta\right)^2+1}\right)\right) d\beta.$$

After solving, we have $L_{\beta}\{\psi\}_{(p)} = -\tan^{-1}(p \ln \beta) + c$ where c is an integral constant.

Taking limit $\beta \to \infty$, we get

$$\lim_{\beta \to \infty} L_{\beta} \{\psi\}_{(p)} = \lim_{\beta \to \infty} (-\tan^{-1}(p \ln \beta) + c)$$
$$0 = -\frac{\pi}{2} + c$$
$$c = \frac{\pi}{2}.$$

And we get

Again, taking limit as $\beta \rightarrow e$

$$\lim_{\beta \to e} L_{\beta} \{\psi\}_{(p)} = \lim_{\beta \to e} \tan^{-1} \left(\frac{1}{p \ln \beta}\right)$$
$$L \left\{\psi(t) = \frac{\sin t}{t}\right\}_{(p)} = \tan^{-1} \left(\frac{1}{p}\right).$$

Example 2. The Laplace integral transform $\frac{e^{-\frac{1}{t}}}{t^{\frac{3}{2}}}, t > 0$. By the

definition of β -Laplace transform

$$L_{\beta}\{\psi\}_{(p)} = \int_0^\infty \beta^{-pu} \psi(u) du.$$

Taking $\psi(t) = \frac{e^{-\frac{1}{t}}}{e^{\frac{3}{2}}}, t > 0$, we have

$$L_{\beta}\{\psi\}_{(p)} = \int_0^\infty \beta^{-pu} \frac{e^{-\frac{1}{u}}}{u^{\frac{3}{2}}} du.$$

Taking partial derivative with respect to $\boldsymbol{\beta}$ (Partial derivative under integral sign) We obtain

$$\begin{split} \frac{\partial}{\partial \beta} \left(L_{\beta} \{\psi\}_{(p)} \right) &= \frac{\partial}{\partial \beta} \int_{0}^{\infty} \beta^{-pu} \frac{e^{-\frac{1}{u}}}{u^{\frac{3}{2}}} du = \int_{0}^{\infty} \left(-\frac{pu}{\beta} \right) \! \beta^{-pu} \frac{e^{-\frac{1}{u}}}{u^{\frac{3}{2}}} du \\ &= - \left(\frac{p}{\beta} \right) \! \int_{0}^{\infty} \beta^{-pu} \frac{e^{-\frac{1}{u}}}{u^{\frac{1}{2}}} du \\ &= - \left(\frac{p}{\beta} \right) \! \sqrt{\frac{\pi}{p \ln \beta}} e^{-2\sqrt{p \ln \beta}}. \end{split}$$

Substitute $p \ln \beta = v$, $\frac{p}{\beta} d\beta = dv$ and integrating with respect to β

$$\begin{split} & L_{\left(\frac{v}{e^{p}}\right)}\{\psi\}_{(p)} = -\int \left(\sqrt{\frac{\pi}{v}}e^{-2\sqrt{v}}\right) dv \\ & L_{\left(\frac{v}{e^{p}}\right)}\{\psi\}_{(p)} = \sqrt{\pi} \cdot e^{-2\sqrt{v}} + c \\ & L_{\beta}\{\psi\}_{(p)} = \sqrt{\pi} \cdot e^{-2\sqrt{p\ln\beta} + c}. \end{split}$$

Taking $\beta \rightarrow \infty$, we have

$$\lim_{\beta \to \infty} L_{\beta} \{\psi\}_{(p)} = \lim_{\beta \to \infty} (\sqrt{\pi} \cdot e^{-2\sqrt{p \ln \beta}} + c)$$
$$0 = 0 + c \Longrightarrow c = 0$$
$$L_{\beta} \{\psi\}_{(p)} = \sqrt{\pi} \cdot e^{-2\sqrt{p \ln \beta}}.$$

Taking limit $\beta \rightarrow e$, we get Laplace transform of the given function

$$L_{\beta}\{\psi\}_{(p)} = \sqrt{\pi} \cdot e^{-2\sqrt{p}}.$$

5. Conclusion

It is obvious that if we apply β -Laplace integral transform via Leibnitz rule, then it is applicable to find Laplace integral transform of many of the functions without including much more calculation.

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